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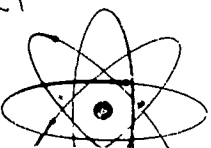


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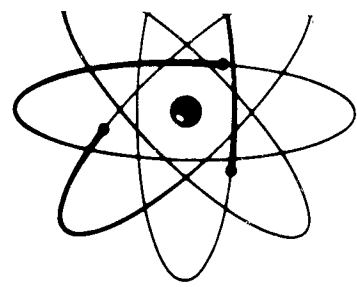


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GAMMA-GAMMA ANGULAR CORRELATION
IN THE DECAY OF COBALT-60

by
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Submitted in partial fulfillment of
the requirements for the degree of
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IN
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This work is accepted as fulfilling
the thesis requirements for the degree of
MASTER OF SCIENCE

IN
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ABSTRACT

This is a continuation of the work of assembling and testing the equipment necessary to conduct gamma-gamma angular correlation experiments, started by Fort A. Vesser, Jr. at the US Naval Postgraduate School (4).

No substantial changes were made in the equipment, but rather a rerun of the Cobalt-60 experiment was made to confirm the results obtained in the previous work. A comparison of the previous work confirms that there is a definite distortion in the shape of the observed correlation function.

The statistical analysis of the data and the solid angle corrections required for evaluation of the results have been programmed for calculation by use of the 1604 Digital Computer. This paper outlines these programs in detail.

The author wishes to express his appreciation for the assistance given him by Professor Harry E. Handley and for his guidance throughout this investigation. Appreciation is expressed to Professor Edmund A. Milne for his advice and comments. Captain T. R. Abernathy, USMC, was very helpful in providing valuable assistance and discussions on various aspects of the computer programming. The cooperation and assistance of Mr. Mervyn C. Brillhart in maintaining and testing the electronic equipment was indispensable.

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CHAPTER I INTRODUCTION

The decay of Cobalt-60 to an excited state of Nickel-60 and the angular correlation of the resulting gamma-gamma cascade have been studied extensively experimentally (2,6,8,9, and others) and the results of the correlation measurements agree very well with the theoretical correlation for the decay scheme shown in Figure 1. Consequently, the decay of Cobalt-60 is now used as one standard for the testing of correlation equipment.

Angular correlation equipment was assembled at the US Naval Postgraduate School by Versar in 1960 for the investigation of the properties of low-lying excited nuclear states. His testing of the operating characteristics of the equipment by the angular correlation of the cascading gamma rays resulting from the decay of Cobalt-60 yielded an observed correlation curve whose shape was not in agreement with the theoretical curve. This work was initiated in order to attempt to verify Versar's results.

In order to facilitate the statistical analysis of the data, a program has been worked out for the use of the 1604 Digital Computer. The program follows the procedures outlined by M.E. Rose (2) and is contained completely in Appendix I.

An additional program has been written to compute the solid angle corrections, which must be applied to the theoretical curve, by the method of M.E. Rose (2). This program is outlined in Appendix II.

The results of this work agreed within statistical expectation with those obtained by Versar, indicating that the unexpected shape of the curve is apparently real and not an instrumental effect. These results are given in Chapter VI.

PARTIAL DECAY SCHEME OF COBALT-60

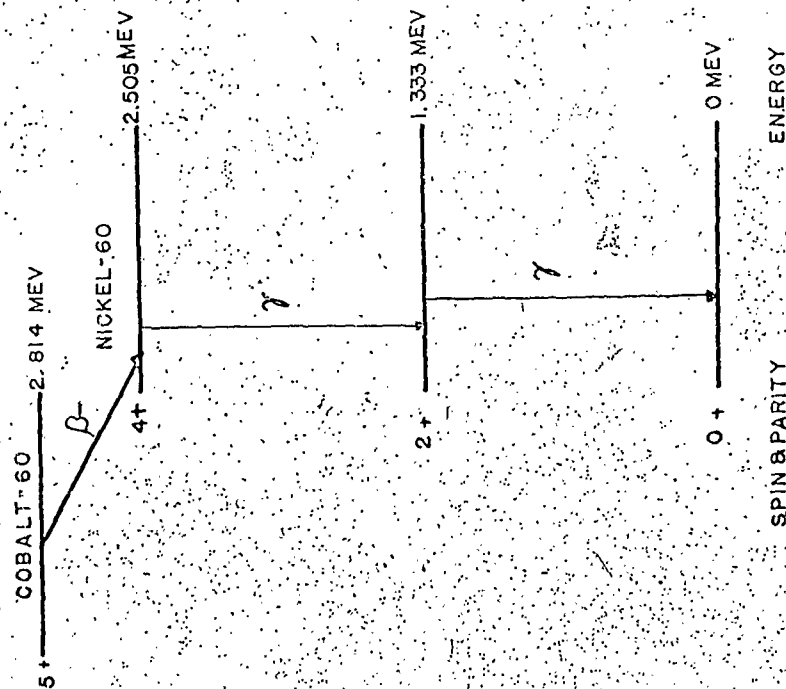


FIGURE 1 2

CHAPTER II THEORETICAL CONSIDERATIONS

The determination of the angular momentum quantum numbers, or spins, of low-lying short-lived excited nuclear states by the nuclear spectroscopic tool, known as angular correlation of successively emitted radiations, is quite common practice today. Only within recent years, because of the vast advancement of electronics, has practical utilization of this tool become possible.

Since nuclei, under ordinary conditions, are randomly oriented in space, it is not possible to observe a radiation pattern. This is due to the fact that the probability of emission of a radiation by an excited nucleus, depends on the angle between the nuclear spin axis and the direction of the emission. Therefore, in order to observe a radiation pattern, it is first necessary to orient the nuclei.

One method of orienting certain nuclei consists of placing the sample in an electric field gradient or a strong magnetic field at a very low temperature. Another method, the one used in the present consideration, consists of selecting nuclei whose nuclear states have spins lying in preferred directions. This happens to be the case for the intermediate state if a nucleus decays by successive emission of two radiations.

In an angular correlation experiment the first radiation is observed in a fixed direction. This establishes an axis to which the direction of emission of the second radiation, originating from the state formed by the first, can be referred. The second radiation has preferred angles of emission with respect to the direction of the first.

In order that the correlation exhibit maximum anisotropy, the angular momentum vector of the intermediate state must not change

direction significantly before the second radiation occurs. The mean life of the intermediate state must be less than about 10^{-8} seconds, the typical precession period of the nuclear spin about local perturbing fields. The angular correlation function is determined by the nature of the radiations and by the spins of the states involved, and thus its measurement may lead to an unambiguous choice of spins for the states.

The theoretical expressions for the correlation functions have been worked out for a large number of cases of interest (7). For the particular case of a gamma-gamma cascade, the correlation function $W(\theta)$ is

$$W(\theta) = \sum_{n=0}^{n_{\max}} A_n P_n(\cos \theta) = \sum_{n=0}^{n_{\max}} P_n \cos^n \theta \quad (1)$$

where P_n is the Legendre Polynomial of even order n , n_{\max} is the smallest even integer of the set of three numbers consisting of twice the multipole order of each of the two gamma rays, and θ is the angle between the directions of emission of the two gamma rays.

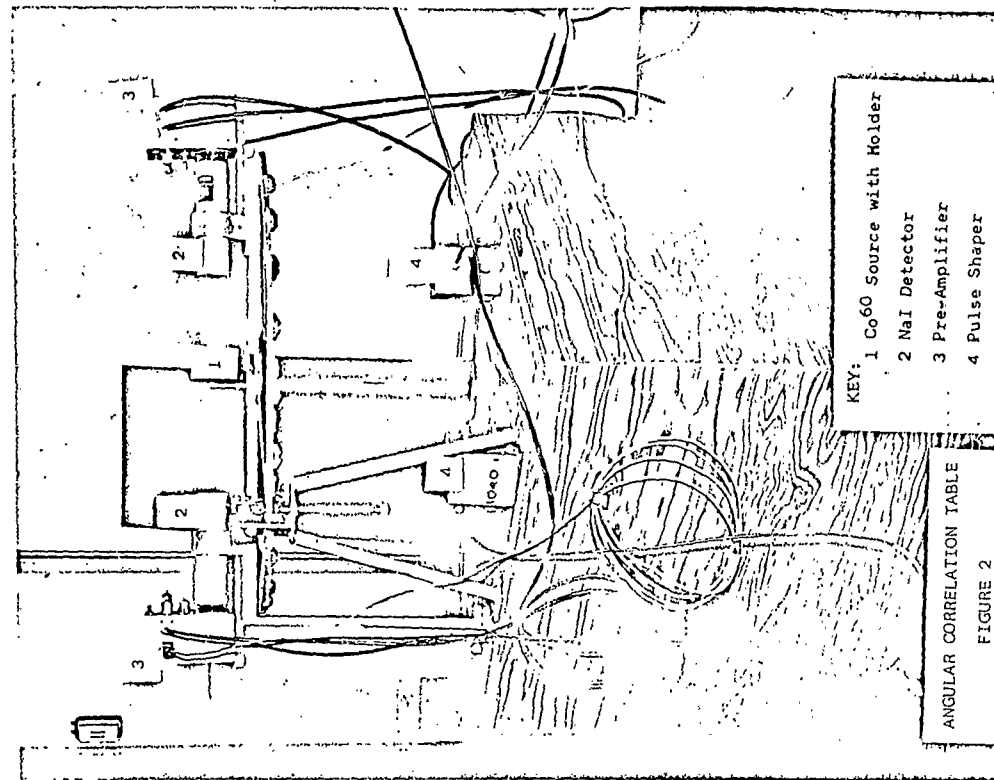
The theoretical correlation function for the cascade in Nickel-60, assuming point detectors and a point source, is

$$W(\theta) = 1 + 0.1020 P_2(\cos \theta) + 0.0091 P_4(\cos \theta) \quad (2)$$

with an anisotropy

$$R = \frac{W(180^\circ) - W(90^\circ)}{W(90^\circ)} = 0.1667 \quad (3)$$

Many workers, after correcting these theoretical results for the effect of non-point detectors have found good agreement between them and their experimental correlation results.



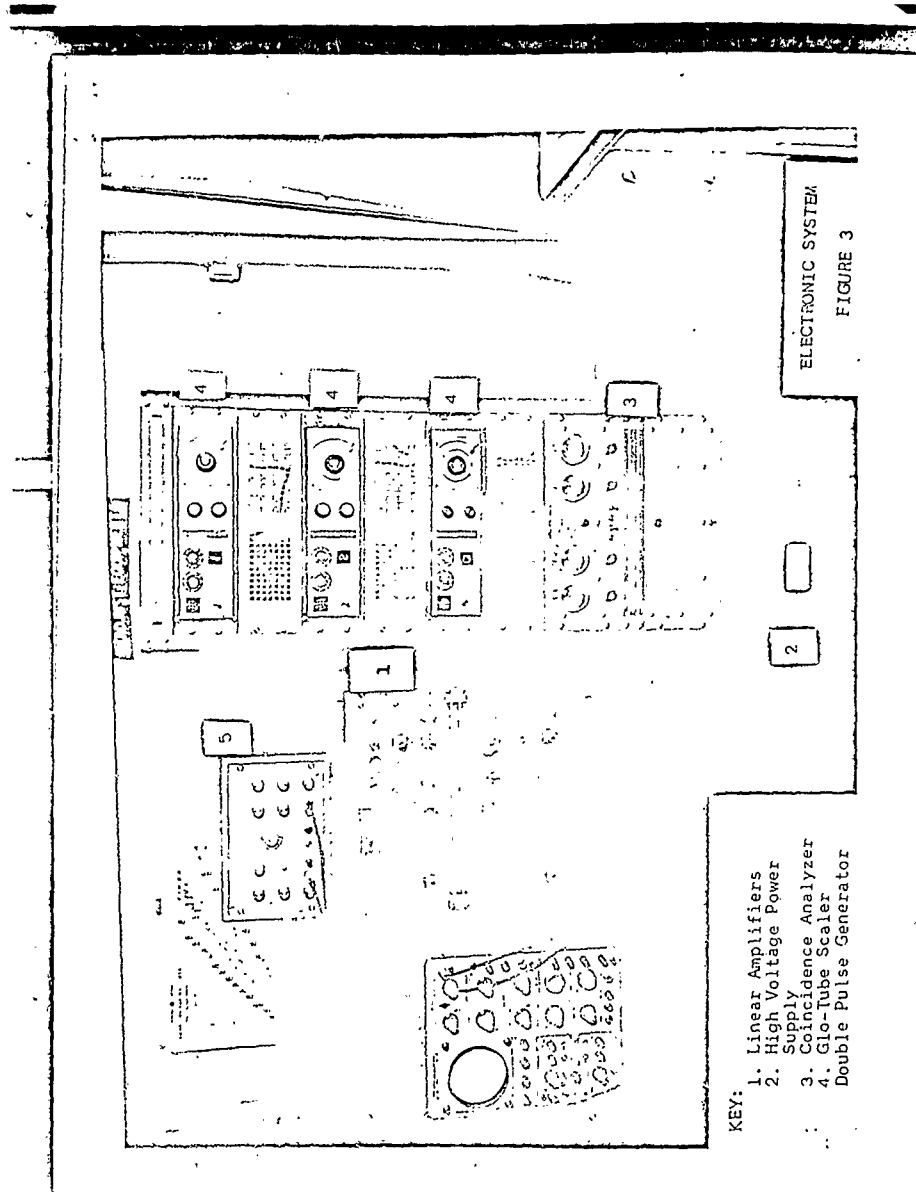
CHAPTER III DESCRIPTION OF EQUIPMENT

The equipment needed for the gamma-gamma correlation measurements must furnish means for detecting the radiation, pulse-height amplification with selection or discrimination, coincidence analysis, and adequate recording of the counts in the various counting channels. A mechanical system must also be provided for mounting the detectors and changing their relative positions and for holding and positioning the radioactive source.

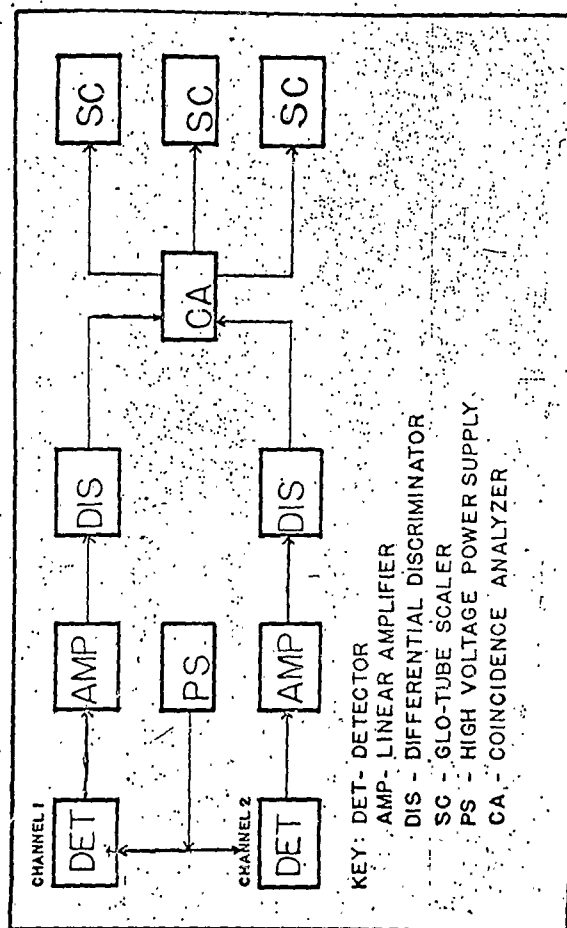
The apparatus used for this work was that assembled by Verser and described in detail in his report. Therefore, only a brief description of the equipment and of the minor alterations performed on it will be presented.

Figure 2 is a photograph of the mechanical setup. The original base of the correlation table has been replaced, the new base is of solid brass with a bearing mounted shaft to enable accurate manual detector positioning.

The electronic equipment (Figure 3) has the functions of detection of the radiation, pulse-height amplification and discrimination, coincidence analysis, and recording the number of counts. The only change in the electronic equipment was in the pulse-shaping networks used in the initial calibration of the apparatus. The pulse shapers were rebuilt to give a better simulation between the test pulses and the actual pulses encountered under operating conditions.



- KEY:
- 1. Linear Amplifiers
 - 2. High Voltage Power Supply
 - 3. Coincidence Analyzer
 - 4. Glo-Tube Scaler
 - 5. Double Pulse Generator



SCHEMATIC DIAGRAM OF THE ELECTRONIC SYSTEM

FIGURE 4

CHAPTER IV EXPERIMENTAL ANALYSIS

The angular correlation function $W(\theta)$ (Equation 1) must be expressed in terms of experimentally-observable quantities. The equipment records the total number of coincidences as a function of the angle between the detectors, and the total number of radiations captured in each channel. C_t , the total number of coincidences in a counting period, is given by

$$C_t = C_g + C_a + C_b \quad (4)$$

where

C_g = the number of genuine coincidences due to the two cascading radiations,
 C_a = the number of accidental coincidences due to the finite resolving time of the electronic equipment, and
 C_b = the number of background coincidences.

The half-life of the source is long compared to the time required to collect the necessary data; therefore, equation (4) can be rearranged and rewritten in terms of counting rates:

$$\dot{C}_g = \dot{C}_t - \dot{C}_a - \dot{C}_b \quad (5)$$

\dot{C}_g can also be expressed in terms of other parameters:

$$\dot{C}_g = \langle \epsilon_1 \alpha_{11} \epsilon_2 \alpha_{22} \epsilon_1 \alpha_{12} \epsilon_2 \alpha_{21} \epsilon_1 \alpha_{11} \epsilon_2 \alpha_{22} \rangle \dot{N}_0 \bar{W}(\theta) \quad (6)$$

where

\dot{N}_0 = the absolute source strength
 θ_{jk} = the fraction of nuclei decaying by the emission of gamma ray "j" and gamma ray "k" in a cascade

A_i = the fractional solid angle subtended by the crystal in Channel "i"
 ρ_{ij} = the fraction of gamma rays "j" which leave the source in the right direction to enter crystal "i" that are not absorbed before reaching the crystal
 ϵ_{ij} = the total efficiency of the crystal in Channel "i" for gamma ray "j"
 α_{ij} = the fraction of the number of gamma rays "j" captured by the crystal in Channel "i" whose pulse heights are accepted by the energy discrimination of Channel "i"
 $\bar{W}(\theta_m)$ = the angular correlation function at an angle θ_m averaged over the solid angles used.

\dot{N}_i , the counting rate in Channel "i", can also be expressed as a function of certain of these same parameters.

$$\dot{N}_i = A_i \dot{N}_0 \sum_j \epsilon_{ij} \alpha_{ij} \rho_{ij} + \dot{N}_{ib} \quad (7)$$

where

\dot{N}_i = the number of gamma rays "i" per disintegration
 \dot{N}_{ib} = the background counting rate in channel "i"

From Equations (5) and (6) we get for the correlation function,

$$\bar{W}(\theta_m) = \frac{\dot{C}_g}{\dot{N}_0} = \frac{\dot{C}_t - \dot{C}_a - \dot{C}_b}{\dot{N}_0} \quad (8)$$

where f is the appropriate function of the ϵ 's, α 's, ρ 's, and θ 's.

By combining Equations (7) and (8), we get for the correlation function

$$\bar{W}(\theta_m) = \frac{(C_t - C_a - C_b) N_0 g}{(N_1 - N_{1b})(N_2 - N_{2b})} \quad (9)$$

where g is a function of the ϵ 's, α 's, β 's, λ 's, δ 's.

In this form $\bar{W}(\theta_m)$ is relatively insensitive to equipment stability. Therefore, it is only necessary to calculate the modified correlation function given by

$$\bar{W}'(\theta_m) = \frac{\bar{W}(\theta_m)}{g_0} = \frac{(C_t - C_a - C_b)}{(N_1 - N_{1b})(N_2 - N_{2b})} \quad (10)$$

The correlation runs were done with the movable detector at seven angles: 90° , 105° , 120° , 135° , 150° , 165° , and 180° . A minimum of six two-hour counting periods were completed at each angle, and the angular order in which they were taken was varied to minimize any systematic fluctuations. Each run yielded single channel counting rates of approximately 1200 counts per second and a total coincidence rate of approximately 1.1 counts per second.

Before the experimental results can be compared to theory, the data must be checked and corrected, if necessary, for the various deviations from the ideal arrangement, i.e., centered point source, no spurious coincidences due to the coincidence resolving time, no scattering, no background, and perfect stability of the electronic equipment. In addition, the theoretical correlation function must be corrected for non-zero solid angle detectors. The methods for accounting for these effects are discussed in the following paragraphs.

A. SOURCE SIZE AND POSITIONING

As a test of the positioning of the source in the center of the correlation table, the integral counting rate of a portion of the Cobalt-60 spectrum was measured in each channel as a function of the angle between the counters. This was done by Verser and checked again after the base of the correlation table had been changed. Within statistical expectations, the counting rates were constant. The source has a volume of about 0.001 cubic inches and can be regarded as a point source at the distance at which the detectors were placed from it.

B. SCATTERING

Compton scattering occurring outside the detectors can give rise to unwanted coincidences. These in general will tend to smear out the measured correlation function. This problem is minimized by use of scintillation crystals as detectors and accepting only the full-energy peaks of the desired gamma rays. The method for determining the pulse height selection limits has been discussed in detail by Verser.

C. BACKGROUND CORRECTION

With the discriminators at the same settings as were used during the correlation runs, background runs were made before and after the correlation runs. The background coincidence counting rate was negligibly small, and the background rate of each channel was such as to yield total background counts of the order of 0.4 percent of the total counts in each channel during the correlation runs. These background rates were used as corrections to the gross counting rates obtained during correlation runs.

D. DETERMINATION OF GENUINE COINCIDENCE RATE

Since the background coincidence counting rate proved experimentally

to be negligible, Equation (5) becomes

$$\dot{C}_g = \dot{C}_t - \dot{C}_a$$

For the total counting time, T , the total number of coincidences is $\dot{C}_t T$ and the standard deviation in the total coincidence rate is

$$\sigma_t = \sqrt{\frac{\dot{C}_t}{T}} \quad (11)$$

Therefore, the standard deviation of the genuine coincidence rate is

$$\sigma_g^2 = \sigma_t^2 + \sigma_a^2 = \frac{\dot{C}_t}{T} + \frac{\dot{C}_a}{T} \quad (12)$$

The accidental coincidence rate is given by

$$\dot{C}_a = 2 \dot{C}_t \dot{N}_1 \dot{N}_2 \sim 2 \dot{C}_t^2 \dot{N}_0 \quad (13)$$

where

\dot{C}_t = Coincidence resolving time of the coincidence analyzer.

In order to minimize the accidental coincidence rate, the source strength must be kept as small as possible. However, if the source strength is decreased, the counting time must be increased to obtain a specified precision in the counting rate. Thus, a compromise must be reached between precision and counting time.

The smallest resolving time obtainable with the present equipment is experimentally determined to be

$$\dot{C}_t = (0.1820 \pm 0.0016) \text{ microseconds;}$$

this resulted in an accidental coincidence rate that was about 30% of the

total coincidence rate. The approximate source strength of the water solution of CoCl_2 used in this experiment was 0.05 millicuries of Cobalt-60.

E. REFERENCE COUNTING RATE

With the discriminator in each channel adjusted to accept only the full-energy peaks of the desired gamma rays, the integral counting rate of the Cobalt-60 spectrum in each channel was determined. This gave a reference counting rate for each channel. During the course of the correlation runs, either the discriminator settings or the gains of the linear amplifiers were adjusted before each run so that, the integral counting rate agreed within ± 5 cps with these reference rates (Table 1).

TABLE 1
REFERENCE COUNTING RATES

$$\begin{aligned} N_1 &= (1085.26 \pm 0.57) \text{ cps} \\ N_2 &= (1117.32 \pm 0.59) \text{ cps} \end{aligned}$$

As a measure of the gain fluctuations during each counting run, the integral counting rate of each channel was measured after the run using the same discriminator setting that was employed during the run. The difference between this counting rate and the reference counting for that channel was determined. If the variation exceeded ± 20 cps in either channel, the run was discarded.

This procedure resulted in about 30% of the runs being discarded because they fell outside of the maximum acceptable fluctuation. This indicates that the electronic system is not as stable as it should be. A check of the line voltage revealed substantial fluctuations during any 24-hour period. The temperature of the room was not controlled in any manner and varied considerably. These could have been the major causes of the large counting rate fluctuations since the discriminators are highly sensitive to both temperature and voltage variations.

P. SOLID ANGLE CORRECTIONS

The solid angle correction which must be applied to the theoretical correlation function to account for the non-zero solid angle of the detectors and enable a comparison to be made with the measured function is described by Rose (2). The correction for two cylindrical crystals whose axes intersect at the source involves the numerical evaluation of integrals of the form

$$I_{n,i} = \int_0^{\gamma_i} P_n(\cos \theta) \{1 - e^{-\gamma_i X_i(\theta)}\} \sin \theta d\theta \quad (14)$$

where

γ_i = the full-energy absorption coefficient of the detector in Channel "i"; For the gamma ray;

γ_i = the half-angle subtended by the front face of the crystal in Channel "i";

$$X_i(\theta) = \tau_i \sec \theta \text{ for } 0 \leq \theta \leq \tan^{-1} \frac{\tau_i}{h_i + \tau_i};$$

$$X_i(\theta) = \tau_i \sec \theta - h_i \sec \theta \text{ for } \tan^{-1} \frac{\tau_i}{h_i + \tau_i} \leq \theta \leq \gamma_i;$$

where

h_i = the distance from the source to the crystal in Channel "i";

τ_i = the thickness of the crystal in Channel "i";

r_i = the radius of the crystal in Channel "i".

The attenuation factors, Q_n , are functions of these integrals:

$$Q_n = \left[\frac{\gamma_{n,1}}{\gamma_{0,1}} \right] \left[\frac{\gamma_{n,2}}{\gamma_{0,2}} \right] \quad (15)$$

The corrected correlation function becomes

$$\bar{W}(\theta) = \sum_{n=0}^{N_{\max}} Q_n A_n P_n(\cos \theta)$$

CHAPTER V EXPERIMENTAL RESULTS

The modified correlation function for the k^{th} run at angle θ_m given by Equation (10) is

$$\bar{W}'_k(\theta_m) = \frac{\left\{ \frac{C_c - C_a - C_b}{(N_1 - N_b)(N_2 - N_b)} \right\}_k}{\left\{ \frac{C_c - C_a - C_b}{(N_1 - N_b)(N_2 - N_b)} \right\}_k}$$

This means that the seven counting rates C_c , C_a , C_b , N_1 , N_{1b} , N_2 , and N_{2b} must be determined for each run at each angle and the standard deviation of each rate calculated.

The background counting rates N_{1b} and N_{2b} for each run were experimentally determined. The single-channel counting rates corrected for scaling losses yielded N_1 and N_2 . The accidental counting rates were calculated for each run from Equation (13). The total coincidence counting rate, C_c , was experimentally determined. The background coincidence rate, C_b , was negligible throughout the experiment.

The experimental value of the modified correlation function $\bar{W}'_k(\theta_m)$ and its standard deviation were calculated for each run at each angle. From these the weighted mean $\bar{W}'(\theta_m)$ and its standard deviation were computed for each angle.

The weighted mean of the measurements taken at each angle was determined by using the inverse square of the standard deviation of each measurement as a weighting factor:

$$\bar{W}'(\theta_m) = \frac{\sum_{k=1}^{N_m} \frac{\bar{W}'_k(\theta_m)}{\sigma_k^2}}{\sum_{k=1}^{N_m} \frac{1}{\sigma_k^2}} \quad (16)$$

where N_m = number of runs at the angle θ_m .

The standard deviation of the weighted mean (\bar{W}) was calculated from the individual standard deviations:

$$\frac{1}{\sigma^2(\bar{W})} = \sum_{i=1}^{M_m} \frac{1}{\sigma_i^2} \quad (19)$$

A comparison between the standard deviations of the means calculated from Equation (19) and those calculated from

$$\sigma^2(\bar{W}) = \sum_{i=1}^{M_m} \frac{[W(\theta_i) - \bar{W}(\theta)]^2}{N_m(N_m - 1)} \quad (20)$$

revealed no significant differences. Therefore the standard deviation obtained from Equation (19) were used in subsequent calculations and the method discussed in paragraph E Chapter IV for treating gain instabilities is realistic.

The values of the weighted means and their standard deviations are tabulated in Table 2.

TABLE 2
The Weighted Means $\bar{W}(\theta_m)$

θ_m (degrees)	Weighted Mean	Standard Deviation
90	5046	45
105	4925	33
120	5109	21
135	5211	36
150	5556	35
165	5667	24
180	5907	43

A curve of the form $\bar{W}(\theta) = \sum_{n=0}^4 A_n P_n(\cos \theta)$ was fitted to the experimental $\bar{W}(\theta_m)$ by the method of least squares as outlined by Rose (2). This analysis was programmed to be carried out by the 1604 computer. (See Appendix II for details).

The values of A_n are given by the matrix equation

$$\|A\| = \frac{1}{\|A^T A\|} \cdot \frac{1}{\|A^T\|} \cdot \frac{1}{\|W\|} \cdot \frac{1}{\|\mu\|} \quad (21)$$

where

$\|A\|$ = the 7 x 3 matrix consisting of the Legendre Polynomials, P_0, P_1, P_2, P_3, P_4 , evaluated at the seven angles (Figure 5A).

$\|A^T\|$ = the transposed A matrix

$\|W\|$ = a diagonal matrix with elements consisting of the inverse squares of the standard deviations of the experimental

$\bar{W}(\theta_m)$ (Figure 5B)

$\|\mu\|$ = a one column matrix whose elements are the experimental values of $\bar{W}(\theta_m)$ (Figure 5C)

$\|A\|$ = the one column matrix whose elements are respectively the least-square coefficients A_0, A_1, A_2 and A_4 (Figure 5D)

$\|A^T A\|^{-1}$ = the inverse of the matrix obtained by multiplying the three designated matrices together. The diagonal terms of this matrix are the squares of the standard deviations of the least-square coefficients, A_n .

The experimental values for the coefficients A_n and their standard deviations, normalized so that $A_0 = 1$, are listed in Table 3. The corresponding values from Verser's work are included for comparison.

TABLE 3
Least Square Coefficients

Verter	This Work
A_0^1	1.0000 ± 0.0034
A_2^1	0.0961 ± 0.0055
A_4^1	0.0339 ± 0.0071

MATRIX CONFIGURATIONS

$P_0(\cos 90^\circ)$	$P_2(\cos 90^\circ)$	$P_4(\cos 90^\circ)$
$P_0(\cos 105^\circ)$	$P_2(\cos 105^\circ)$	$P_4(\cos 105^\circ)$
$P_0(\cos 120^\circ)$	$P_2(\cos 120^\circ)$	$P_4(\cos 120^\circ)$
$P_0(\cos 135^\circ)$	$P_2(\cos 135^\circ)$	$P_4(\cos 135^\circ)$
$P_0(\cos 150^\circ)$	$P_2(\cos 150^\circ)$	$P_4(\cos 150^\circ)$
$P_0(\cos 165^\circ)$	$P_2(\cos 165^\circ)$	$P_4(\cos 165^\circ)$
$P_0(\cos 180^\circ)$	$P_2(\cos 180^\circ)$	$P_4(\cos 180^\circ)$

$\sigma^{-2}(90^\circ)$	0	0	0	0	0
$\sigma^{-2}(105^\circ)$	0	0	0	0	0
$\sigma^{-2}(120^\circ)$	0	$\sigma^{-2}(120^\circ)$	0	0	0
$\sigma^{-2}(135^\circ)$	0	0	$\sigma^{-2}(135^\circ)$	0	0
$\sigma^{-2}(150^\circ)$	0	0	0	$\sigma^{-2}(150^\circ)$	0
$\sigma^{-2}(165^\circ)$	0	0	0	0	$\sigma^{-2}(165^\circ)$
$\sigma^{-2}(180^\circ)$	0	0	0	0	$\sigma^{-2}(180^\circ)$

(A)

$\bar{W}^1(90^\circ)$	A_0^1
$\bar{W}^1(105^\circ)$	A_2^1
$\bar{W}^1(120^\circ)$	A_4^1
$\bar{W}^1(135^\circ)$	
$\bar{W}^1(150^\circ)$	
$\bar{W}^1(165^\circ)$	
$\bar{W}^1(180^\circ)$	

(C)

A_0^0
A_2^0
A_4^0

(D)

FIGURE 5

20

A. Experimental Anisotropies

The anisotropy is defined:

$$\bar{R} = \frac{\bar{W}(180^\circ) - \bar{W}(90^\circ)}{\bar{W}(90^\circ)} \quad (22)$$

The anisotropy for the least-square curve was calculated from the experimental correlation function. Since the anisotropy can be determined with much greater precision than can the least-square coefficients themselves, it provides a more sensitive means for comparing the experimental and theoretical functions. The experimental anisotropy is

$$\bar{R} = 0.167 \pm 0.013$$

The standard deviation was calculated by the formula presented by Klena and McGowan (6). The computer program for the calculation of \bar{R} and its standard deviation is outlined in Appendix I.

CHAPTER VI RESULTS AND CONCLUSIONS

A. COMPARISON BETWEEN EXPERIMENT AND THEORY

The experimental and corrected theoretical functions for this experiment are:

Experimental:

$$\bar{W}(\theta) = (1.0000 \pm 0.0026) + (0.0971 \pm 0.0057) P_2 + (0.0231 \pm 0.0066) P_4$$

Theoretical:

$$\bar{W}(\theta) = 1.0000 \pm 0.0996 P_2 + 0.00838 P_4$$

The experimental anisotropy was calculated to be $\bar{R} = 0.167 \pm 0.013$ and is compared with Verser's value of $\bar{R} = 0.171 \pm 0.027$ and with the theoretical anisotropy corrected for the solid angle, $\bar{R} = 0.161$.

Figure 6 shows the experimental results and their standard deviations for this work and that of Verser. Also shown are the least-square correlation curves for both sets of data and the corrected theoretical correlation curve. All the curves have been normalized to unity at $\theta = 90^\circ$.

B. DISCUSSION AND CONCLUSIONS

From Figure 6, it can be seen that the curve obtained in this work agrees with that obtained by Verser. Both of these curves show an apparently real dip in the region from 105° to 135° . It is only a remote possibility that this dip can be attributed to statistical fluctuations, and it is unlikely to be due to instrumental effects.

Two conjectures of possible mechanisms are:

- (a) The source may have crystallized resulting in preferential alignment of the nuclei within the crystal due to internal electric fields.

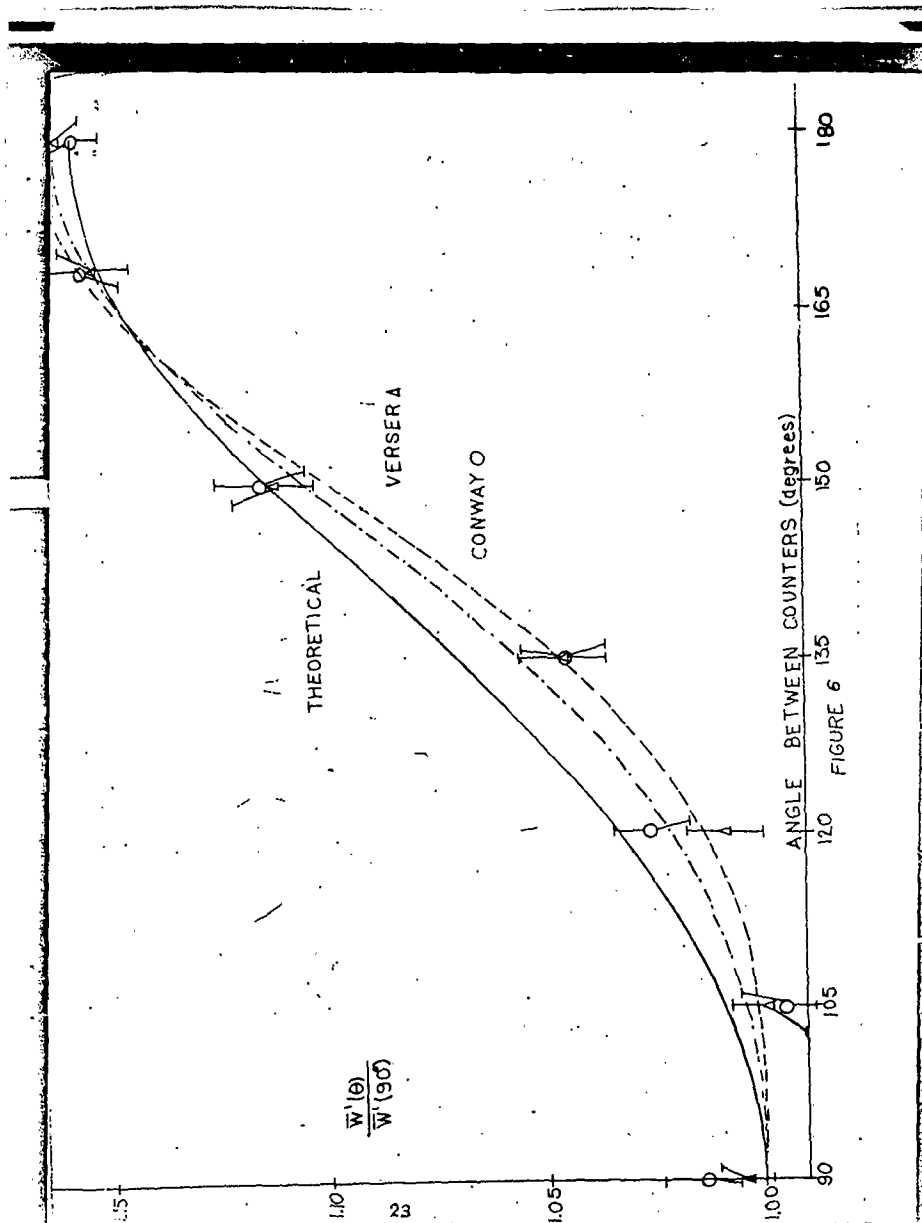


FIGURE 6

- (b) The ferromagnetic properties of Cobalt could have led to a preferential alignment of the nuclei in the earth's magnetic field.

The following series of experiments are suggested for future work.

- (a) The results for measurements conducted for angles between 180° and 270° should be compared with those already obtained.
- (b) In order to check on the possible effects of orientation due to crystalline electric fields, the experiment should be repeated with the present source at various orientations relative to the fixed counter.
- (c) In order to check on the possibility of orientation by magnetic effects the source should be placed in a magnetic field of known direction and the measurements repeated.
- (d) A correlation curve using a powdered Cobalt-60 source in place of the present source should be obtained for comparison.

These experiments would determine conclusively if any effects causing preferential nuclear orientation are present.

The following instrumentation improvements are suggested in order to improve the stability of the equipment.

- (a) A very stable power supply should be purchased to aid in preventing the line voltage fluctuations from interfering with the electronic system.
- (b) The effect of temperature variations on the electronic system can be minimized by placing the equipment in a temperature controlled room for conducting future work.
- (c) A slow coincidence unit whose resolving time is independent of the counting rate should be obtained.

The precision of the results should be improved by the addition of a fast coincidence channel working in conjunction with the slow coincidence unit to reduce the number of accidental coincidences.

Since the experimental curves using Cobalt-60 have not agreed with the theoretical curve, the experiments on Cobalt-60 should be continued until the source of the distortion is ascertained. Perhaps the effect can be related to the physical properties of the CoCl_2 source.

Least-Square Fit of Experimental Results APPENDIX I

The comparison of the experimental correlation function to the functional form of the theoretical correlation function given by Equation (1) is accomplished by the method of least-squares, as outlined by Rose (2). The analysis requires the solution of Equation (20) for the values of A'_n .

The computer solution of the matrix equation, like the solid angle calculation, is carried out in floating point format. In addition to the solution for the values of A'_n , the program also calculates the anisotropy by Equation (21), and its standard deviation by the formula presented by Klemm and McGowan (6). A flow sheet outline of the program is given in Figure (7). In order to provide a certain amount of flexibility to the program, the angles used have been left as a parameter to be supplied. The program as written is, however, limited to seven different angular positions. The parameters needed for the program may be placed anywhere in the computer and their addresses supplied to the program by placing them in the B registers in the following order:

B1 = weighted means of the correlation function in ascending angular order.

B2 = the standard deviations of the weighted means in the same order as above.

B3 = the values of the angles used in ascending order.

The program is available for future use on both biocatal tape in machine language, and IBM cards in assembly language. The biocatal tape contains the following:

1. The main program for calculation of A'_n
2. The program of normalization of the A'_n
3. The program for calculation of the resulting correlation function
4. The normalization of the correlation function
5. The calculation of the anisotropy

LEAST-SQUARE FIT FLOW CHART

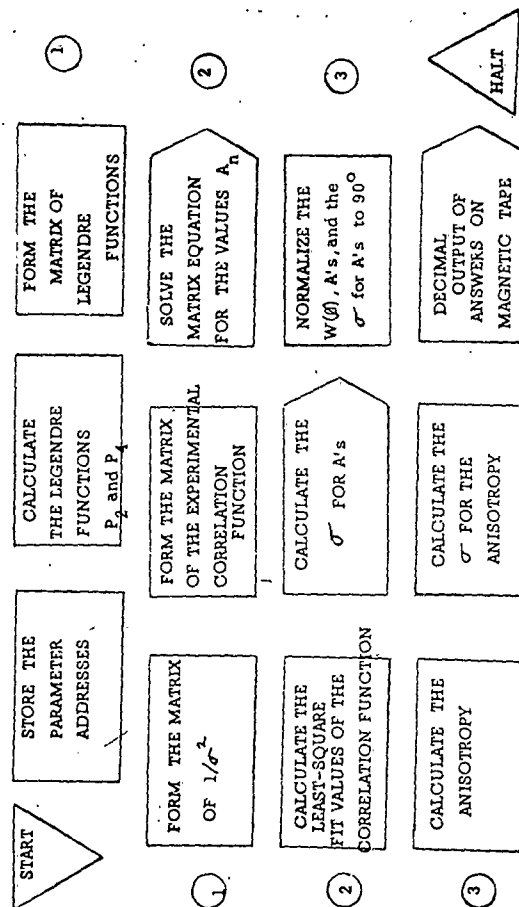


FIGURE 7

6. The calculation of the standard deviation for the anisotropy

7. The subroutines

- a. Floating Point Sine-Cosine
- b. Floating Point Square Root
- c. Floating Point Decimal Output.

The results of the computations are dumped on magnetic tape unit (4) in the following sequence:

Lines 1 thru 3 The matrix inverse

Lines 4 thru 6 Check answer; included in the program is a check of the inverse which consists of multiplying the inverse by the original matrix

Line 7 The values of A'_0 , A'_2 and A'_4

Line 8 The standard deviations of the A'_n

Line 9 thru 13 The resulting correlation functions

Line 14 and 15 The normalized correlation functions

Line 16 The normalized A'_n

Line 17 The normalized standard deviations of the A'_n

Line 18 The anisotropy

Line 19 The square of the standard deviation of the anisotropy

STATISTICAL ANALYSIS PROGRAM

```

CNG 12000
REM BURTON J CONWAY APRIL 1961
REM LEAST SQUARE FIT OF ANGULAR
REM CORRELATION DATA.
REM MUST BE SUPPLIED WITH
REM PARAMETER ADDRESSES AS FOLLOWS.
REM B1 CORRELATION FUNCTION.
REM B2 SIGMA OF CORRELATION FUNCT. ON.
REM B3 ANGLES USED IN DEGREES.
EQU 13000
ZEUS
EQU 13100
RHO
EQU 13200
TRANS
EQU 13250
ATW
EQU 13260
ATWA
EQU 13500
INVER
EQU 13520
CHECK
EQU 13540
LAT
EQU 13570
LATW
EQU 13620
ANS
EQU 13650
PTNIS
EQU 60000
SINFL
EQU 10000
MATRIX

```


12013	20 0 12461	STA 0 DON
	32 0 12320	FMU 0 FLT2
12014	31 0 12316	FSB 0 FLT
	33 0 12317	FDV 0 FLT1
12015	20 1 12326	STA 1 WORK2
	12 0 12461	LDA 0 DON
12016	32 0 12322	FMU 0 FLT4
	20 0 12462	STA 0 DON+1
12017	12 0 12461	LDA 0 DON
	32 0 12461	FMU 0 DON
12020	32 0 12323	FMU 0 FLT5
	31 0 12462	FSB 0 DON+1
12021	30 0 12320	FAD 0 FLT2
	33 0 12321	FDV 0 FLT3
12022	20 1 12335	STA-1 WORK4
	50 0 00000	ENI 0 0
12023	54 1 00006	ISK 1 6
	75 0 12012	SLJ 0 NEXT
12024	12 0 12325	LDA 0 WORK1
	20 2 12434	STA 2 ABLE
12025	51 2 00001	INI 2 1
	12 1 12326	LDA 1 WORK2
12026	20 2 12434	STA 2 ABLE
	51 2 00001	INI 2 1

12000	56 1 12045	ANISO	EQV	13624
	50 1 00000	STU	1 FORMU	
12001	57 2 12032	ENI	1 0	
	56 2 12033	SIL	2 SIGMA	
12002	50 2 00000	STU	2 SIGMA+1	
	50 0 00000	ENI	2 0	
12003	56 3 12004	ENI	0 0	
	50 3 00000	STU	3 /+1	
12004	12 1 00000	ENI	3 0	
	32 0 12512	LDA	1 0	
12005	75 4 60001	FMU	0 CONV	
	50 0 00000	SLJ	4 SINEL+1	
12006	50 0 00000	ENI	0 0	
	50 0 00000	ENI	0 0	
12007	50 0 00000	ENI	0 0	
	50 0 00000	ENI	0 0	
12010	20 1 12463	STA	1 TEMP	
	50 0 00000	ENI	0 0	
12011	54 1 00006	ISK	1 6	
	75 0 12004	SLJ	0 START	
12012	12 1 12463	LDA	1 TEMP	
	32 1 12463	FMU	1 TEMP	

STARTS FORMATION
OF THE LEG. POL.

12027	12 1 12335	LDA 1 MGRK4
	20 2 12434	STA 2 ABLE
12030	51 2 00001	INI 2 1
	50 0 00000	ENI 0 0
12031	54 1 00006	ISK 1 6
	75 0 12024	SLJ 0 FORM
12032	12 0 12316	SIGMA
	33 1 00000	FDV 1 0
12033	33 1 00000	FDV 1 0
	20 1 12463	STA 1 TEMP
12034	54 1 00006	ISK 1 6
	75 0 12032	SLJ 0 SIGMA
12035	50 2 00000	ENI 2 0
	50 1 00000	ENI 1 0
12036	12 1 12333	FORMCAD LDA-1 CAD
	22 0 12043	ATP 0 /+5
12037	12 2 12463	LDA 2 TEMP
	20 1 12333	STA 1 CAD
12040	51 1 00001	INI 1 1
	50 0 00000	ENI 0 0
12041	54 2 00006	ISK 2 6
	75 0 12036	SLJ 0 FORMCAD

12042	75 0 12044	SLJ 0 /+2
	50 0 00000	ENI 0 0
12043	51 1 00001	INI 1 1
	75 0 12036	SLJ 0 FORMCAD
12044	50 1 00000	ENI 1 0
	50 2 00000	ENI 2 0
12045	12 1 00000	FORMU LDA 1 0
	20 1 12344	STA 1 MU
12046	54 1 00006	ISK 1 6
	75 0 12045	SLJ 0 FORMU
12047	50 1 00000	ENI 1 0
	50 2 00000	ENI 2 0
12050	75 4 10000	SLJ 4 MATRIX
	50 0 00000	ENI 0 0
12051	50 0 00000	ENI 0 0
	50 0 00000	ENI 0 0
12052	50 0 00000	ENI 0 0
	50 0 00002	ENI 0 2
12053	50 0 00007	ENI 0 7
	50 0 12434	ENI 0 ABLE
12054	50 0 00003	ENI 0 3
	50 0 06626	ENI 0 6626
12055	50 0 00000	ENI 0 0
	50 0 13200	ENI 0 TRANS

TRANSPOSE

12071	50 0 00003	ENI 0 3	
	50 0 13260	ENI 0 ATWA	
12072	75 4 10000	SLJ 4 MATRIX	INVERSE
	50 0 00000	ENI 0 0	
12073	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12074	50 0 00000	ENI 0 0	
	50 0 00005	ENI 0 5	
12075	50 0 00003	ENI 0 3	
	50 0 13260	ENI 0 ATWA	
12076	50 0 00044	ENI 0 360	
	50 0 13272	ENI 0 ATWA+12	
12077	50 0 00000	ENI 0 0	
	50 0 13500	ENI 0 INVER	
12100	75 4 10000	SLJ 4 MATRIX	CHECK
	50 0 00000	ENI 0 0	
12101	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12102	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12103	50 0 00003	ENI 0 3	
	50 0 13500	ENI 0 INVER	
12104	50 0 00003	ENI 0 3	
	50 0 13260	ENI 0 ATWA	

12056	75 4 10000	SLJ 4 MATRIX	ATWA
	50 0 00000	ENI 0 0	
12057	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12060	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12061	50 0 00003	ENI 0 3	
	50 0 13200	ENI 0 TRANS	
12062	50 0 00007	ENI 0 7	
	50 0 12353	ENI 0 CAD	
12063	50 0 00007	ENI 0 7	
	50 0 13230	ENI 0 ATW	
12064	75 4 10000	SLJ 4 MATRIX	ATWA
	50 0 00000	ENI 0 0	
12065	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12066	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12067	50 0 00003	ENI 0 3	
	50 0 13230	ENI 0 ATW	
12070	50 0 00007	ENI 0 7	
	50 0 12434	ENI 0 ABLE	

12105	50 0 00003	ENI 0 3	
	50 0 13520	ENI 0 CHECK	
12106	75 4 10000	SLJ 4 MATRIX	LAT
	50 0 00000	ENI 0 0	
12107	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12110	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12111	50 0 00003	ENI 0 3	
	50 0 13500	ENI 0 INVER	
12112	50 0 00003	ENI 0 3	
	50 0 13200	ENI 0 TRANS	
12113	50 0 00007	ENI 0 7	
	50 0 13540	ENI 0 LAT	
12114	75 4 10000	SLJ 4 MATRIX	LATW
	50 0 00000	ENI 0 0	
12115	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12116	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12117	50 0 00003	ENI 0 3	
	50 0 13540	ENI 0 LAT	

12120	50 0 00007	ENI 0 7	
	50 0 12553	ENI 0 CAD	
12121	50 0 00007	ENI 0 7	
	50 0 13570	ENI 0 LATW	
12122	75 4 10000	SLJ 4 MATRIX	ANS
	50 0 00000	ENI 0 0	
12123	50 0 00000	ENI 0 0	
	50 0 00000	ENI 0 0	
12124	50 0 00000	ENI 0 0	
	50 0 00004	ENI 0 4	
12125	50 0 00003	ENI 0 3	
	50 0 13570	ENI 0 LATW	
12126	50 0 00007	ENI 0 7	
	50 0 12344	ENI 0 MU	
12127	50 0 00001	ENI 0 1	
	50 0 13620	ENI 0 ANS	
12130	50 4 00000	ENI 4 0	
	50 5 00000	ENI 5 0	
12131	50 6 00000	ENI 6 0	
	50 1 00000	ENI 1 0	
12132	50 2 00000	ENI 2 0	
	50 3 00000	ENI 3 0	
12133	54 1 77777	ISK 1 77777	
	75 0 12133	SLJ 0 /	

12134	50 0 00000	ENI 0 0
12135	74 7 42040	EXP 7 42040
12136	75 4 71000	SLJ 4 71000
12137	50 0 00000	ENI 0 0
12138	01 0 13500	01 0 INVER
12139	04 0 00001	04 0 1
12140	12 0 12315	LDA 0 MIKE
12141	70 0 12136	RAD 0 /-1
12142	54 4 00002	ISK 4 2
12143	75 0 12135	SLJ 0 /-3
12144	54 1 77777	ISK 1 77777
12145	75 0 12141	SLJ 0 /
12146	54 1 77777	ISK 1 77777
12147	75 0 12142	SLJ 0 /
12148	50 0 00000	ENI 0 0
12149	74 7 42040	EXP 7 42040
12150	75 4 71000	SLJ 4 71000
12151	50 0 00000	ENI 0 0
12152	01 0 13520	01 0 CHECK
12153	04 0 00004	04 0 4
12154	12 0 12315	LDA 0 MIKE
12155	70 0 12145	RAD 0 /-1

12147	54 4 00002	ISK 4 2
12148	75 0 12144	SLJ 0 /-3
12149	54 1 77777	ISK 1 77777
12150	75 0 12150	SLJ 0 /
12151	50 0 00000	ENI 0 0
12152	74 7 42040	EXP 7 42040
12153	75 4 71000	SLJ 4 71000
12154	50 0 00000	ENI 0 0
12155	01 0 13620	01 0 ANS
12156	04 0 00007	04 0 7
12157	12 0 13620	REM STARTS FORMATION OF THE
12158	32 4 12434	REM CORRELATION FUNCTION
12159	20 0 13100	LDA 0 ANS
12160	51 4 00001	PMU 4 ABLE
12161	12 0 13621	STA 0 REO
12162	32 4 12434	INI 4 1
12163	30 0 13100	LDA 0 ANS+1
12164	20 0 13100	PMU 4 ABLE
12165	12 0 13622	PAD 0 REO
12166	51 4 00001	STA 0 REO
12167	32 4 12434	LDA 0 ANS+2
12168	30 0 13100	INI 4 1
12169	20 0 13100	PMU 4 ABLE
12170	12 0 13622	AMP4 IN A REG
12171	51 4 00001	AMP4 IN A REG
12172	32 4 12434	SUM IN A REG
12173	30 0 13100	SUM IN A REG

ANS TO STORAGE

12162	20 6 13000	STA 6 ZEUS
	51 6 00003	INI 6 3
12163	54 4 00024	ISK 4 24
	75 0 12154	SLJ 0 DIMP
12164	12 2 13500	IDA 2 INVER
	75 4 12527	SLJ 4 SQREL
12165	50 0 00000	ENI 0 0
	50 0 00000	ENI 0 0
12166	20 1 12501	STA 1 TTA
	51 2 00004	INI 2 4
12167	54 1 00002	ISK 1 2
	75 0 12164	SLJ 0 /-3
12170	50 1 00000	ENI 1 0
	50 2 00000	ENI 2 0
12171	50 3 00000	ENI 3 0
	50 4 00000	ENI 4 0
12172	54 1 77777	ISK 1 77777
	75 0 12172	SLJ 0 /
12173	50 0 00000	ENI 0 0
	74 7 42040	EXP 7 42040
12174	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0

12175	01 0 12501	01 0 TTA
	04 0 00010	04 0 10
12176	54 1 77777	ISK 1 77777
	75 0 12176	SLJ 0 /
12177	50 0 00000	ENI 0 0
	74 7 42040	EXP 7 42040
12200	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0
12201	01 0 13000	01 0 ZEUS
	04 0 00011	04 0 11
12202	12 0 12315	LDA 0 MIKE
	70 0 12201	RAD 0 /-1
12203	54 5 00004	ISK 5 4
	75 0 12200	SLJ 0 /-3
12204	12 1 13000	LDA 1 ZEUS
	33 0 13000	FDV 0 ZEUS
12205	20 2 12472	STA 2 NORM
	51 1 00003	INI 1 3
12206	54 2 00006	ISK 2 6
	75 0 12204	SLJ 0 /-2
12207	54 1 77777	ISK 1 77777
	75 0 12207	SLJ 0 /
12210	50 0 00000	ENI 0 0
	74 7 42040	EXP 7 42040

NORMALIZES THE
CORRELATION
FUNCTION

DUMPS THE
NORMALIZED
CORRELATION
FUNCTION

12211	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0
12212	01 0 12472	01 0 NORM
	04 0 00016	04 0 16
12213	12 0 12315	IDA 0 MIKE
	70 0 12212	RAD 0 /-1
12214	54 2 00001	ISK 2 1
	75 0 12211	SLJ 0 /-3
12215	50 1 00000	ENI 1 0
	50 0 00000	ENI 0 0
12216	12 1 13620	IDA 1 ANS
	33 0 13620	FDV 0 ANS
12217	20 1 12507	STA 1 AN
	50 0 00000	ENI 0 0
12220	54 1 00002	ISK 1 2
	75 0 12216	SLJ 0 /-2
12221	54 1 77777	ISK 1 77777
	75 0 12221	SLJ 0 /
12222	50 0 00000	ENI 0 0
	74 7 42040	EXF 7 42040
12223	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0

12240	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0
12241	01 0 13624	01 0 ANISO
	04 0 00022	04 0 22
12242	12 0 12314	IDA 0 ZERO
	20 0 12514	STA 0 ROB
12243	12 1 12344	IDA 1 MU
	31 2 13000	FSB 2 ZEUS
12244	20 0 12513	STA 0 ERSIL
	32 0 12513	FMU 0 ERSIL
12245	32 3 12353	FMU 3 CAD
	30 0 12514	FAD 0 ROB
12246	20 0 12514	STA 0 ROB
	51 2 00003	INI 2 3
12247	51 3 00010	INI 3 10
	50 0 00003	ENI 0 0
12250	54 1 00006	ISK 1 6
	75 0 12243	SLJ 0 AVG
12251	12 0 12514	IDA 0 ROB
	33 0 12324	FDV 0 FLNG
12252	20 0 12514	STA 0 ROB
	50 2 00000	ENI 2 0

12253	50 3 00000	ENI 3 0
	50 1 00000	ENI 1 0
12254	12 0 12514	LDA 0 ROB
	32 1 13500	FMD 1 INVER
12255	20 2 12515	STA 2 AXE
	51 1 00004	INT 1 4
12256	54 2 00002	ISK 2 2
	75 0 12254	SLJ 0 /-2
12257	12 0 12514	LDA 0 ROB
	32 0 13501	FMD 0 INVER+1
12260	20 0 12520	STA 0 COW2
	12 0 12514	LDA 0 ROB
12261	32 0 13505	FMD 0 INVER+5
	20 0 12521	STA 0 COW24
12262	12 0 13621	LDA 0 ANS+1
	30 0 13622	FAD 0 ANS+2
12263	20 0 12522	STA 0 A24
	32 0 12522	FMD 0 A24
12264	20 0 12523	STA 0 A242
	12 0 12522	LDA 0 A24
12265	33 0 13620	FDV 0 ANS
	20 0 12524	STA 0 GET
12266	32 0 12524	FMD 0 GET
	20 0 12524	STA 0 GET

NORMALIZES THE
SIGMAS FOR A

12224	01 0 12507	01 0 AN
	04 0 00020	04 0 20
12225	12 1 12501	LDA 1 TTA
	33 0 13000	FDV 0 ZEUS
12226	20 1 12504	STA 1 TTAN
	50 0 00000	ENI 0 0
12227	54 1 00002	ISK 1 2
	75 0 12225	SLJ 0 /-2
12230	54 1 77777	ISK 1 77777
	75 0 12230	SLJ 0 /
12231	50 0 00000	ENI 0 0
	74 7 42040	EXF 7 42040
12232	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0
12233	01 0 12504	01 0 TTAN
	04 0 00021	04 0 21
12234	12 0 13022	LDA 0 ZEUS+22
	31 0 13000	FSB 0 ZEUS
12235	33 0 13000	FDV 0 ZEUS
	20 0 13624	STA 0 ANISO
12236	54 1 77777	ISK 1 77777
	75 0 12236	SLJ 0 /
12237	50 0 00000	ENI 0 0
	74 7 42040	EXF 7 42040

12267	12 0 12522	LDA 0 A24	12302	32 0 12524	FMJ 0 GET
	32 0 13620	FMJ 0 ANS		20 0 13630	STA 0 FINIS
12270	20 0 12522	STA 0 A24	12303	54 1 77777	ISK 1 77777
	50 0 00000	ENI 0 0		75 0 12303	SLJ 0 /
12271	12 0 12521	LDA 0 CONV24	12304	50 0 00000	ENI 0 0
	32 0 12317	FMJ 0 FLTL		74 7 42040	EXP 7 42040
12272	30 0 12516	FAD 0 AXE+1	12305	75 4 71000	SLJ 4 71000
	30 0 12517	FAD 0 AXE+2		50 0 00000	ENI 0 0
12273	33 0 12523	FDV 0 A242	12306--	01 0 13630	01 0 FINIS
	20 0 12525	STA 0 WIT	12307	04 0 00025	04 0 25
12274	12 0 12515	LDA 0 AXE		54 1 77777	ISK 1 77777
	33 0 13620	FDV 0 ANS	12310	75 0 12307	SLJ 0 /
12275	33 0 13620	FDV 0 ANS		54 1 77777	ISK 1 77777
	30 0 12525	FAD 0 WIT	12311	75 0 12310	SLJ 0 /
12276	20 0 12525	STA 0 WIT		50 0 00000	ENI 0 0
	12 0 12520	LDA 0 CONV2	12312	74 7 42040	EXP 7 42040
12277	30 0 12521	FAD 0 CONV24		74 0 42003	EXP 0 42003
	33 0 12522	FDV 0 A24	12313	50 0 00000	ENI 0 0
12300	32 0 12317	FMJ 0 FLTL		76 0 12000	SLS 0 12000
	20 0 12526	STA 0 WIT+1	12314	50 0 00000	ENI 0 0
12301	12 0 12525	LDA 0 WIT		00 0 00000	OCT 0
	31 0 12526	FSB 0 WIT+1	12315	00 0 00004	MIKE 0 0 4
				00 0 00001	0 0 1

12316	20 0 14000	FLT	DEC	1.0
	00 0 00000			
12317	20 0 24000	FLT1	DEC	2.0
	00 0 00000			
12320	20 0 26000	FLT2	DEC	3.0
	00 0 00000			
12321	20 0 44000	FLT3	DEC	8.0
	00 0 00000			
12322	20 0 57400	FLT4	DEC	30.0
	00 0 00000			
12323	20 0 64300	FLT5	DEC	35.0
	00 0 00000			
12324	20 0 34000	FLT6	DEC	4.0
	00 0 00000			
12325	20 0 14000	WORK1	DEC	1.0
	00 0 00000			
12326	00 0 00000	WORK2	BSS	7
	00 0 00000			
12335	00 0 00000	WORK4	BSS	7
	00 0 00000			
12344	17 5 34167	MU	DEC	.5046D-6
	17 3 10477			

12345	17 5 34103	DEC	.4925D-6
	21 2 75007		
12346	17 5 34222	DEC	.5109D-6
	23 0 76631		
12347	17 5 34276	DEC	.5211D-6
	06 7 01135		
12350	17 5 34522	DEC	.5556D-6
	22 1 24417		
12351	17 5 34601	DEC	.5667D-6
	75 2 36515		
12352	17 5 34751	DEC	.5907D-6
	02 3 02276		
12353	20 7 07231	CAD	.657462D-17
	17 2 43106		
12354	00 0 00000	OCT	0
	00 0 00000		
12355	00 0 00000	OCT	0
	00 0 00000		
12356	00 0 00000	OCT	0
	00 0 00000		
12357	00 0 00000	OCT	0
	00 0 00000		
12360	00 0 00000	OCT	0
	00 0 00000		

12374	00 0 00000	OCT 0
12375	00 0 00000	OCT 0
12376	00 0 00000	OCT 0
12377	00 0 00000	OCT 0
12400	00 0 00000	OCT 0
12401	00 0 00000	OCT 0
12402	00 0 00000	OCT 0
12403	20 7 05176 27 1 03164	DEC -4725897D+17
12404	00 0 00000	OCT 0
12405	00 0 00000	OCT 0
12406	00 0 00000	OCT 0
12407	00 0 00000	OCT 0

12361	00 0 00000	OCT 0
12362	00 0 00000	OCT 0
12363	20 7 05367 05 2 01171	DEC -4938271D+17
12364	00 0 00000	OCT 0
12365	00 0 00000	OCT 0
12366	00 0 00000	OCT 0
12367	00 0 00000	OCT 0
12370	00 0 00000	OCT 0
12371	00 0 00000	OCT 0
12372	00 0 00000	OCT 0
12373	20 7 05015 20 7 37137	DEC -4526935D+17

12423	20 7 15041	DEC	.9124087D+17
12424	16 1 31732	OCT	0
12425	00 0 00000	OCT	0
12426	00 0 00000	OCT	0
12427	00 0 00000	OCT	0
12428	00 0 00000	OCT	0
12429	00 0 00000	OCT	0
12430	00 0 00000	OCT	0
12431	00 0 00000	OCT	0
12432	00 0 00000	OCT	0
12433	20 7 05367	DEC	.4938271D+17
12434	05 2 01171	DEC	1.0
12435	20 0 14000 / ABLE	DEC	1.0
12436	00 0 00000	DEC	1.0
12437	57 7 73777	DEC	1.0
12438	77 7 77777	DEC	1.0
12439	17 7 66000	DEC	1.0
12440	00 0 00000	DEC	1.0

12410	00 0 00000	OCT	0
12411	00 0 00000	OCT	0
12412	00 0 00000	OCT	0
12413	20 7 04477	DEC	.4164951D+17
12414	37 2 35173	OCT	0
12415	00 0 00000	OCT	0
12416	00 0 00000	OCT	0
12417	00 0 00000	OCT	0
12418	00 0 00000	OCT	0
12419	00 0 00000	OCT	0
12420	00 0 00000	OCT	0
12421	00 0 00000	OCT	0
12422	00 0 00000	OCT	0
12423	00 0 00000	OCT	0
12424	00 0 00000	OCT	0
12425	00 0 00000	OCT	0
12426	00 0 00000	OCT	0
12427	00 0 00000	OCT	0
12428	00 0 00000	OCT	0
12429	00 0 00000	OCT	0
12430	00 0 00000	OCT	0
12431	00 0 00000	OCT	0
12432	00 0 00000	OCT	0
12433	00 0 00000	OCT	0
12434	00 0 00000	OCT	0
12435	00 0 00000	OCT	0
12436	00 0 00000	OCT	0
12437	00 0 00000	OCT	0
12438	00 0 00000	OCT	0
12439	00 0 00000	OCT	0
12440	00 0 00000	OCT	0

12452	17 7 26004	DEC	.023471
12453	31 0 11317	DEC	1.0
12454	20 0 14000	DEC	.8995
12455	00 0 00000	DEC	.684614
12456	20 0 07144	DEC	1.0
12457	26 4 16254	DEC	1.0
12458	20 0 05364	DEC	1.0
12459	13 3 47506	DEC	1.0
12460	20 0 14000	DEC	1.0
12461	00 0 00000	DEC	1.0
12462	00 0 00000	DEC	1.0
12463	00 0 00000	DEC	1.0
12464	00 0 00000	DEC	1.0
12465	00 0 00000	DEC	1.0
12466	00 0 00000	DEC	1.0
12467	00 0 00000	DEC	1.0
12468	00 0 00000	DEC	1.0
12469	00 0 00000	DEC	1.0
12470	00 0 00000	DEC	1.0
12471	00 0 00000	DEC	1.0
12472	00 0 00000	DEC	1.0
12473	00 0 00000	DEC	1.0
12474	00 0 00000	DEC	1.0
12475	00 0 00000	DEC	1.0
12476	00 0 00000	DEC	1.0
12477	00 0 00000	DEC	1.0
12478	00 0 00000	DEC	1.0
12479	00 0 00000	DEC	1.0
12480	00 0 00000	DEC	1.0
12481	00 0 00000	DEC	1.0
12482	00 0 00000	DEC	1.0
12483	00 0 00000	DEC	1.0
12484	00 0 00000	DEC	1.0
12485	00 0 00000	DEC	1.0
12486	00 0 00000	DEC	1.0
12487	00 0 00000	DEC	1.0
12488	00 0 00000	DEC	1.0
12489	00 0 00000	DEC	1.0
12490	00 0 00000	DEC	1.0
12491	00 0 00000	DEC	1.0
12492	00 0 00000	DEC	1.0
12493	00 0 00000	DEC	1.0
12494	00 0 00000	DEC	1.0
12495	00 0 00000	DEC	1.0
12496	00 0 00000	DEC	1.0
12497	00 0 00000	DEC	1.0
12498	00 0 00000	DEC	1.0
12499	00 0 00000	DEC	1.0

12437	20 0 14000	DEC	1.0
12438	00 0 00000	DEC	1.0
12439	60 0 11466	DEC	1.0
12440	73 1 10267	DEC	1.0
12441	17 7 54455	DEC	1.0
12442	71 7 14727	DEC	1.0
12443	20 0 14000	DEC	1.0
12444	00 0 00000	DEC	1.0
12445	60 0 23777	DEC	1.0
12446	77 7 77777	DEC	1.0
12447	60 0 13277	DEC	1.0
12448	76 5 40722	DEC	1.0
12449	20 0 14000	DEC	1.0
12450	00 0 00000	DEC	1.0
12451	17 7 61000	DEC	1.0
12452	10 7 23316	DEC	1.0
12453	60 0 11400	DEC	1.0
12454	05 1 74265	DEC	1.0
12455	20 0 14000	DEC	1.0
12456	00 0 00000	DEC	1.0
12457	20 0 05000	DEC	1.0
12458	00 0 00000	DEC	1.0
12459	00 0 00000	DEC	1.0
12460	00 0 00000	DEC	1.0
12461	00 0 00000	DEC	1.0
12462	00 0 00000	DEC	1.0
12463	00 0 00000	DEC	1.0
12464	00 0 00000	DEC	1.0
12465	00 0 00000	DEC	1.0
12466	00 0 00000	DEC	1.0
12467	00 0 00000	DEC	1.0
12468	00 0 00000	DEC	1.0
12469	00 0 00000	DEC	1.0
12470	00 0 00000	DEC	1.0
12471	00 0 00000	DEC	1.0
12472	00 0 00000	DEC	1.0
12473	00 0 00000	DEC	1.0
12474	00 0 00000	DEC	1.0
12475	00 0 00000	DEC	1.0
12476	00 0 00000	DEC	1.0
12477	00 0 00000	DEC	1.0
12478	00 0 00000	DEC	1.0
12479	00 0 00000	DEC	1.0
12480	00 0 00000	DEC	1.0
12481	00 0 00000	DEC	1.0
12482	00 0 00000	DEC	1.0
12483	00 0 00000	DEC	1.0
12484	00 0 00000	DEC	1.0
12485	00 0 00000	DEC	1.0
12486	00 0 00000	DEC	1.0
12487	00 0 00000	DEC	1.0
12488	00 0 00000	DEC	1.0
12489	00 0 00000	DEC	1.0
12490	00 0 00000	DEC	1.0
12491	00 0 00000	DEC	1.0
12492	00 0 00000	DEC	1.0
12493	00 0 00000	DEC	1.0
12494	00 0 00000	DEC	1.0
12495	00 0 00000	DEC	1.0
12496	00 0 00000	DEC	1.0
12497	00 0 00000	DEC	1.0
12498	00 0 00000	DEC	1.0
12499	00 0 00000	DEC	1.0

12507	00 0 00000	AN	BSS	3	REM	FLOATING POINT SQUARE ROOT
	00 0 00000				SLJ	0 0 ALARM EXIT
12512	17 7 24357	CONV	DEC	.01745329252	STU	1 SQFEXIT STORE R1
	50 6 50451				AJP	0 SQFEXIT-2 TEST 0 ARGUMENT
12513	00 0 00000	ERSIL	BSS	1	AJP	3 SQFEL TEST - ARGUMENT
	00 0 00000				ENT	0 0
12514	00 0 00000	ROB	BSS	1	SCM	0 SQMASKI SET CORRECT
	00 0 00000				ALS	0 1 SIGN FOR
12515 -	00 0 00000	AXE	BSS	3	IRS	0 37D EXPONENT
	00 0 00000				AJP	3 /+1 COMPUTE
12520	00 0 00000	COV02	BSS	1	INA	0 2 EXPONENT
	00 0 00000				INA	0 3777 FEE ROOT
12521	00 0 00000	COV24	BSS	1	IRS	0 1 SHIFT EXPONENT
	00 0 00000				ALS	0 36D AND SIGN
12522	00 0 00000	A24	BSS	1	STA	0 SQERAS
	00 0 00000				IDL	0 SQMASK2
12523	00 0 00000	A242	BSS	1	ARS	0 4
	00 0 00000				QTP	3 /+1
12524	00 0 00000	GET	BSS	1	ARS	0 1
	00 0 00000				STA	0 SQERAS+1
12525	00 0 00000	MIT	BSS	2	AID	0 SQFCOVS
	00 0 00000				STA	0 SQERAS+2
					LDA	0 SQFCOVS2

12542	27 0 12565	DVF 0 SQERAS+2
12543	14 0 12553	ADD 0 SQFCOM1
12543	50 1 00001	ENI 1 1
	04 0 00000	ENQ 0 0
12544	20 0 12565	STA 0 SQERAS+2
	12 0 12562	LDA 0 SQERAS+1
12545	27 0 12563	DVF 0 SQERAS+2
	14 0 12563	ADD 0 SQERAS+2
12546	01 0 00001	ARS 0 1
	55 1 12544	LTP 1 /-2
12547	01 0 00011	ARS 0 11
	42 0 12561	SOM 0 SQERAS
12550	32 0 12560	RMI 0 SQ4MSK3
	52 1 12527	LJU 1 SQFPL
12551	51 1 00001	INTL 1 1
	57 1 12552	STL 1 /+1
12552	50 1 00000	SQFEXTM ENI 1 0
	75 0 00000	SLJ 0 0
12553	21 3 60405	SQFCOM1 DEC 218518306
	01 4 47436	
12554	76 3 72106	SQFCOM2 DEC -30228999
	43 6 47567	
12555	03 0 56173	SQFCOM3 DEC 15451577
	52 7 71341	

APPENDIX II

Solid Angle Correction

The solid angle correction, which must be applied to the theoretical correlation function to enable a comparison to be made with the experimental correlation function, is described by Rose (1). The correction involves the numerical evaluation of the integrals given by Equation (14):

$$I_{11} = \int_0^{\gamma} p_n(\cos\theta) (1 - e^{-\tau} X_1(\theta)) \sin\theta d\theta$$

and the subsequent solutions of Equations (15) and (16).

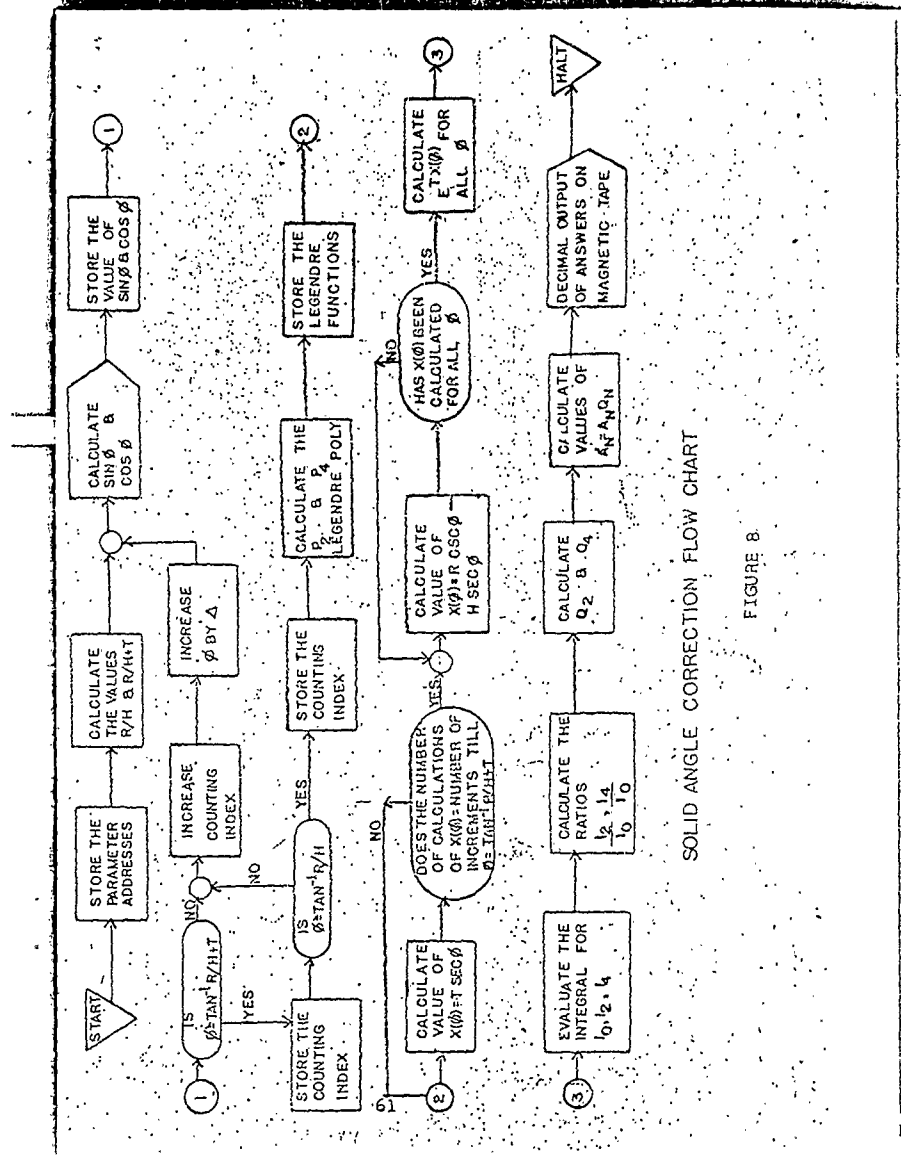
The computer solution for this correction was carried out in floating point format. The numerical integration was computed using an interval of 0.001 radians. A flow sheet outline of the program is given in Figure (8). The program must be supplied with the parameters τ , h , r , and t same in both channels (see Table 4) thus enabling solid angle corrections to be made for any experimental set-up to which these corrections may be applicable. These parameters may be placed anywhere in the computer and their addresses supplied to the program by placing them in the B registers in the following order

- B1 = r (address of the crystal radius in centimeters)
- B2 = t (address of the crystal thickness in centimeters)
- B3 = h (address of the crystal to source distance in centimeters)
- B4 = γ (the value of the full-energy absorption coefficient of the detector for the gamma ray in centimeters⁻¹)

The values of the parameters used in the calculation of the corrected angular correlation coefficients for the theoretical correlation function and the results of the calculations are listed in Table 4.

The program is available for future use on biocatal tape in machine language and on IBM cards in assembly language. The biocatal tape contains all the necessary subroutines to enable a complete computation, i.e.

1. The main program for solution of Equation (14),
2. The subroutines



SOLID ANGLE CORRECTION FLOW CHART

FIGURE 8

- a. Floating Point Sine-Cosine
- b. Floating Point Exponential
- c. Floating Point Decimal Output.

The results of the computations are dumped on magnetic tape unit (4) in the following sequence:

Line 1 I_0
 Line 2 I_2
 Line 3 I_4
 Line 4 I_2/I_0
 Line 5 I_4/I_0
 Line 6 Q_2
 Line 7 A_2Q_2
 Line 8 Q_4
 Line 9 A_4Q_4

TABLE 4
 SOLID ANGLE CORRECTION PARAMETERS

h (cm)	18.0
t (cm)	2.54
z (cm)	2.54
T (cm ⁻²)	0.098
I_0	0.0079
I_2	0.0078
I_4	0.0076
I_2/I_0	0.9882
I_4/I_0	0.9609
Q_2	0.9765
Q_4	0.9234
Q_2A_2	0.0996
Q_4A_4	0.0084

SOLID ANGLE CORRECTION PROGRAM

```

ORG 12000
REM BURTON J COMWAY APRIL 1961
REM SOLID ANGLE CORRECTIONS
REM OF THE THEORETICAL CORRELATION
REM CURVE. MUST BE SUPPLIED
REM WITH PARAMETER ADDRESSES AS
REM FOLLOWS.
REM B1 ADDRESS OF XTAL RADIUS.
REM B2 ADDRESS OF XTAL THICKNESS.
REM B3 ADDRESS OF CRYSTAL TO
REM SOURCE DISTANCE.
REM B4 ADDRESS OF PARAMETER TAU.

HOLD EQU 5500
TEMP EQU 6500
WORK2 EQU 7500
WORK4 EQU 10500
HANS EQU 11500
SOLD EQU 12500
JAY0 EQU 13500
JAY2 EQU 13504
JAY4 EQU 13510
JAY20 EQU 13514
  
```

```

JAY40 EQU 13520
EXPTL EQU 60600
SINFL EQU 60000
12000 LDA 1 0
20 0 12163 STA 0 RAD
12001 LDA 2 0
20 0 12164 STA 0 THK
12002 LDA 3 0
20 0 12165 STA 0 HGT
12003 LDA 4 0
20 0 12166 STA 0 TAU
12004 LDA 0 ZRO
20 5 13500 STA 5 JAY0
12005 ISK 5 30
75 0 12004 SLJ 0 /-1
12006 LDA 0 ZRO
20 0 12151 STA 0 DELTA
12007 ENI 1 0
50 2 00000 ENI 2 0
12010 50 3 00000 ENI 3 0
50 4 00000 ENI 4 0
12011 50 5 00000 ENI 5 0
50 6 00000 ENI 6 0
  
```

12012	12 0 12163	START	LDA 0 RAD	FORMATION OF R/H
	33 0 12165		FDV 0 HGT	
12013	20 0 12172		STA 0 PLEAS	
	50 0 00000		ENI 0 0	
12014	12 0 12165		LDA 0 HGT	FORMATION OF R/H+T
	30 0 12164		FAD 0 THK	
12015	20 0 12175		STA 0 ONCE	
	12 0 12163		LDA 0 RAD	
12016	33 0 12175		FDV 0 ONCE	
	20 0 12175		STA 0 ONCE	
12017	12 0 12151	PGW	LDA 0 DELTA	SINE AND COSINE
	50 0 00000		ENI 0 0	FORMATION
12020	75 4 60001		SLJ 4 SINF+1	COSINE
	50 0 00000		ENI 0 0	
12021	50 0 00000		ENI 0 0	
	50 0 00000		ENI 0 0	
12022	50 0 00000		ENI 0 0	
	50 0 00000		ENI 0 0	
12023	20 1 06500		STA 1 TEMP	COSINE TO STORAGE
	12 0 12151		LDA 0 DELTA	
12024	75 4 60000		SLJ 4 SINF+1	SINE ROUTINE
	50 0 00000		ENI 0 0	
12025	50 0 00000		ENI 0 0	
	50 0 00000		ENI 0 0	

12026	50 0 00000	ENI 0 0		
	50 0 00000	ENI 0 0		
12027	20 1 06500	STA 1 HOLD		SINE TO STORAGE
	33 1 06500	FDV 1 TEMP		SINE/COSINE
12030	31 0 12175	REP+1	FSE 0 ONCE	SINE/COSINE-R/H+T
	22 0 12141		AJP 0 FIRST	JUMPS IF A IS 0
12031	22 2 12141		AJP 2 FIRST	JUMPS IF A IS +
	51 1 00001	INI 1 1		
12032	12 0 12152	LDA 0 DELTA+1		
	30 0 12151	FAD 0 DELTA		
12033	20 0 12151	STA 0 DELTA		
	75 0 12017	SLJ 0 BGN		
12034	56 1 12105	RENT	SIU 1 JAN+3	STARTS FORMATION OF
	56 1 12073		SIU 1 INCL+5	THE LEGENDS VOL.
12035	56 1 12111		SIU 1 JAN+3	
	56 1 12051		SIU 1 RAK	
12036	56 1 12074		SIU 1 INCL+6	
	56 1 12101		SIU 1 JAN+1	
12037	56 1 12115		SIU 1 JAN+3	
	50 1 00000	ENI 1 0		
12040	12 1 06500	LDA 1 TEMP		COSB IN A REG.
	32 1 06500	FAD 1 TEMP		

12041	20 0 12173	STA 0 DON		
	32 0 12157	FMJ 0 FIT2		
12042	31 0 12155	FSB 0 FIT		
	33 0 12156	FDV 0 FIT1		P2 IN A REG.
12043	20 1 07500	STA 1 WORK2		
	12 0 12173	LDA 0 DON		
12044	32 0 12161	FMJ 0 FIT4		
	20 0 12174	STA 0 DON+1		
12045	12 0 12173	LDA 0 DON		
	32 0 12173	FMJ 0 DON		
12046	32 0 12162	FMJ 0 FIT5		
	31 0 12174	FSB 0 DON+1		
12047	30 0 12157	FAD 0 FIT2		P4 IN THE A REG.
	33 0 12160	FDV 0 FIT3		
12050	20 1 10500	STA 1 WORK4		
	50 0 00000	ENI 0 0		
12051	54 1 00000	ENK		NUMBER OF ANGLES
	75 0 12040	SLJ 0 NEXT+4		INSERTED BY PROGRAM
12052	12 0 12164	HALF		FORMATION OF XB
	33 1 06500	FDV 1 TEMP		
12053	20 1 11500	STA 1 HANS		
	50 0 00000	ENI 0 0		
12054	54 1 00000	ISK 1 0		NUMBER INSERTED
	75 0 12052	SLJ 0 HALF		BY PROGRAM

12055	12 1 11500	LDA 1 HANS		
	32 0 12166	FMJ 0 TAU		
12056	20 1 11500	STA 1 HANS		
	50 0 00000	ENI 0 0		NUMBER ENTERED
12057	54 1 00000	ISK 1 0		BY PROGRAM
	75 0 12055	SLJ 0 /-2		STARTS FORMATION
12060	12 1 11500	POWER		OF E TO THE TAU X
	75 4 60600	LDA 1 HANS		
12061	50 0 00000	ENI 0 0		
	50 0 00000	ENI 0 0		
12062	20 1 12500	STA 1 SOLD		
	50 0 00000	ENI 0 0		
12063	54 1 00000	ISK 1 0		
	75 0 12060	SLJ 0 POWER		
12064	54 2 00001	ISK 2 1		
	75 0 12066	SLJ 0 /+2		
12065	75 0 12076	SLJ 0 HANG		
	50 0 00000	ENI 0 0		
12066	50 1 00000	INCL		
	51 1 00001	INI 1 1		
12067	12 0 12165	LDA 0 HGT		STARTS CALCULATION
	33 1 06500	FDV 1 TEMP		R CSCR-SECB

12070	20 1 12500	STA 1 SOLD
	12 0 12163	IDA 0 RAD
12071	33 1 05500	FDV 1 HOLD
	31 1 12500	FSB 1 SOLD
12072	32 0 12166	FMJ 0 TAU
	20 1 11500	STA 1 HANS
12073	54 1 00000	ISK 1 0
	75 0 12057	SLJ 0 /-4
12074	50 1 00000	ENI 1 0
	56 1 12063	STU 1 POWER+3
12075	50 1 00000	ENI 1 0
	75 0 12060	SLJ 0 POWER
12076	12 0 12155	HANG
	33 1 12500	FDV 1 SOLD
12077	20 1 12500	STA 1 SOLD
	12 0 12155	IDA 0 FLT
12100	31 1 12500	FSB 1 SOLD
	20 1 12500	STA 1 SOLD
12101	54 1 00000	ISK 1 0
	75 0 12076	SLJ 0 HANG
12102	12 0 12153	JAY
	32 1 12500	FMJ 1 SOLD
12103	32 1 05500	FMJ 1 HOLD
	32 0 12152	FMJ 0 DELTA+1

STARTS CALCULATION
OF THE INTERVAL

12104	30 0 13500	FAD 0 JAYO
	20 0 13500	STA 0 JAYO
12105	54 1 00000	ISK 1 0
	75 0 12102	SLJ 0 /-3
12106	12 1 07500	JANTO
	32 1 12500	FMJ 1 SOLD
12107	32 1 05500	FMJ 1 HOLD
	32 0 12152	FMJ 0 DELTA+1
12110	30 0 13504	FAD 0 JAY2
	20 0 13504	STA 0 JAY2
12111	54 1 00000	ISK 1 0
	75 0 12106	SLJ 0 /-3
12112	12 1 10500	JAYTR
	32 1 12500	FMJ 1 SOLD
12113	32 1 05500	FMJ 1 HOLD
	32 0 12152	FMJ 0 DELTA+1
12114	30 0 13510	FAD 0 JAY4
	20 0 13510	STA 0 JAY4
12115	54 1 00000	ISK 1 0
	75 0 12112	SLJ 0 /-3
12116	12 0 13504	IDA 0 JAY2
	33 0 13500	FDV 0 JAYO

NUMBER OF TIMES
TILL R/H

NUMBER OF TIMES
TILL R/H

12133	54 5 00010	ISK 5 10
	75 0 12130	SLJ 0 /-3
12134	54 1 77777	ISK 1 77777
	75 0 12134	SLJ 0 /
12135	54 2 77777	ISK 2 77777
	75 0 12135	SLJ 0 /
12136	50 0 00000	ENI 0 0
	74 7 42040	EXF 7 42040
12137	74 0 42003	EXF 0 42003
	50 0 00000	ENI 0 0
12140	76 0 05000	SLS 0 5000
	50 0 00000	ENI 0 0
12141	12 0 12147	FIRST
	20 0 12030	STA 0 REPL
12142	12 0 12150	LDA 0 SEC+1
	20 0 12031	STA 0 REPL+1
12143	56 1 12054	SIU 1 HALF+2
	56 1 12057	SIU 1 HALF+5
12144	56 1 12063	SIU 1 POWER+3
	56 1 12075	SIU 1 INCL+7
12145	56 1 12066	SIU 1 INCL
	51 1 00001	INI 1 1

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12117	20 0 13514	STA 0 JAX20
	12 0 13510	LDA 0 JAX4
12120	33 0 13500	FDV 0 JAX0
	20 0 13520	STA 0 JAX40
12121	12 0 13514	LDA 0 JAX20
	32 0 13514	FMJ 0 JAX20
12122	20 0 13524	STA 0 JAX40+4
	32 0 12167	FMJ 0 THEOR
12123	20 0 13530	STA 0 JAX40+10
	12 0 13520	LDA 0 JAX40
12124	32 0 13520	FMJ 0 JAX40
	20 0 13534	STA 0 JAX40+14
12125	32 0 12170	FMJ 0 THEOR+1
	20 0 13540	STA 0 JAX40+20
12126	34 1 77777	ISK 1 77777
	75 0 12126	SLJ 0 /
12127	50 0 00000	ENI 0 0
	74 7 42040	EXF 7 42040
12130	75 4 71000	SLJ 4 71000
	50 0 00000	ENI 0 0
12131	01 0 13500	01 0 JAX0
	04 0 00001	04 0 1.
12132	12 0 12154	LDA 0 MIKE
	70 0 12131	RAD 0 /-1

72

12146	75 0 12017	SLJ 0 EBN
	50 0 00000	ENI 3 0
12147	51 0 12172	FSB 0 PIFAS
	22 0 12034	AJP 0 NEXT
12150	22 2 12034	AJP 2 NEXT
	51 1 00001	INI 1 1
12151	00 0 00000	DELTA OCT 0
	00 0 00000	
12152	17 6 64061	DEC .001
	11 5 64570	
12153	20 0 14000	WORK DEC 1.0
	00 0 00000	
12154	00 0 00004	MIKE 0 0 4
	00 0 00001	0 0 1
12155	20 0 14000	FLT DEC 1.0
	00 0 00000	
12156	20 0 24000	FLTL DEC 2.0
	00 0 00000	
12157	20 0 26000	FLT2 DEC 3.0
	00 0 00000	
12160	20 0 14000	FLT3 DEC 8.0
	00 0 00000	
12161	20 0 57400	FLTL4 DEC 30.0
	00 0 00000	

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12162	20 0 64300	FLT5 DEC 35.0
	00 0 00000	
12163	20 0 25050	RAD DEC 2.54
	75 3 41217	
12164	20 0 25050	THK DEC 2.54
	75 3 41217	
12165	20 0 54400	HGT DEC 18.0
	00 0 00000	
12166	17 7 46213	TAU DEC .098
	20 7 12601	
12167	17 7 46416	THEOR DEC .1020
	25 4 02030	
12170	17 7 14511	DEC .00907
	51 2 62577	
12171	00 0 00000	ZRO OCT 0
	00 0 00000	
12172	00 0 00000	PIEAS BSS 1
	00 0 00000	
12173	00 0 00000	DON BSS 2
	00 0 00000	
12175	00 0 00000	ONCE BSS 1
	00 0 00000	

END

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